

# Gradient flow formulation and longtime behaviour of a constrained Fokker-Planck equation

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joint work with Simon Eberle and Barbara Niethammer

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# Introduction and Motivation

**Starting point:** The classical Fokker-Planck equation

$$\partial_t \rho(t, x) = \partial_x (\nu^2 \partial_x \rho(t, x) + H'(x) \rho(t, x))$$

with  $\rho(0, x) = \rho^0(x)$  and  $\int \rho^0(x) dx = 1$  posses a free energy

$$\mathcal{F}(\rho) = \nu^2 \int \rho(x) \log \rho(x) dx + \int H(x) \rho(x) dx + \mathcal{F}_0.$$

such that it is the gradient flow with respect to the Wasserstein metric

$$\partial_t \rho = -\mathcal{K}(\rho) D\mathcal{F}(\rho) \quad \text{with} \quad \mathcal{K}(\rho)\varphi := -\partial_x(\rho \partial_x \varphi).$$

Convergence to equilibrium via the entropy method:

Set  $\gamma_0 := \exp(-H/\nu^2)/Z_0$  with  $Z_0 = \int \exp(-H(x)/\nu^2) dx$  and  $\mathcal{F}_0 = \log Z_0$  then

$$\mathcal{F}(\rho) = \nu^2 \mathcal{H}(\rho|\gamma_0) \geq 0 \quad \text{with} = 0 \text{ iff } \rho = \gamma_0.$$

Energy dissipation relation:

$$\frac{d}{dt} \mathcal{F}(\rho(t)) = - \int |\nu^2 \partial_x \log \rho + H'|^2 dx \leq 0.$$

leads to exponential convergence to equilibrium  $\mathcal{F}(\rho(t)) \leq e^{-\frac{t}{\sigma_{\text{LSI}}}} \mathcal{F}(\rho(0))$ .

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### Motivation:

- Many-particle storage system exposed to external dynamical loading obeying

$$\int x \rho(t, x) dx = \ell(t) \quad \text{with } \ell \text{ dynamical load.} \quad \mathcal{C}$$

- Example: Lithium-ion batteries during a charge-discharging cycle<sup>1</sup>

The constraint  $\mathcal{C}$  introduces a **nonlocal Lagrange multiplier**  $\sigma(t)$

$$\partial_t \rho(t, x) = \partial_x \left( \nu^2 \partial_x \rho(t, x) + (H'(x) - \sigma(t)) \rho(t, x) \right),$$

given by  $\sigma(t) = \int H'(x) \rho(t, x) dx + \dot{\ell}(t).$

**First goal:** Incorporate constraint in gradient flow formulation<sup>2</sup>

**Second goal:** Give a characterization of the internal time-scale of relaxation.

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# Constrained gradient flow

## Constrained gradient flow: Formalism

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**Constrained state space:**  $\mathcal{M}^\ell := \{\rho \in \mathcal{P}(\mathbb{R}) : M_1(\rho) = \ell\}$ ,  $M_1(\rho) := \int x\rho(x) dx$ .  
In general  $\mathcal{M}^{\mathcal{C};t} = \{\rho \in \mathcal{P}(\mathbb{R}) : \mathcal{C}(\rho; t) = 0\}$ , here  $\mathcal{C}(\rho; t) := M_1(\rho) - \ell(t)$ .

**Assumption:** Constraint is **nondegenerate**

$$\langle DC(\rho; t), \mathcal{K}(\rho)DC(\rho; t) \rangle = |\nabla \mathcal{C}(\rho; t)|_\rho^2 > 0.$$

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- For stationary constraints  $\mathcal{C}(\rho, t) = \mathcal{C}(\rho, 0)$  an orthogonal projection is sufficient
- Dynamical constraints: augment state space with time dimension  
⇒ pass from nonautonomous system to autonomous

### Constrained gradient flow

$\rho : \mathbb{R}_+ \rightarrow \mathcal{M}$  is the constrained gradient flow with respect to the nondegenerate dynamical constraint  $\mathcal{C} : \mathcal{M} \times \mathbb{R}_+ \rightarrow \mathbb{R}$ , if for all  $t \geq 0$ :

$$\partial_t \rho = -\mathcal{K}(\rho(t))DF(\rho(t)) + \sigma(\rho(t), t) \mathcal{K}(\rho(t))DC(\rho(t), t),$$

with 
$$\sigma(\rho, t) = \frac{\langle DF(\rho), \mathcal{K}(\rho)DC(\rho, t) \rangle + \partial_t \mathcal{C}(\rho, t)}{|\nabla \mathcal{C}(\rho, t)|_\rho^2}.$$

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Let  $h > 0$  be a fixed time step and define for  $k \geq 1$

$$\rho^k = \arg \min_{\rho \in \mathcal{M}^{\ell(kh)}} \left( \frac{1}{2} W_2^2(\rho^{k-1}, \rho) + h\mathcal{F}(\rho) \right).$$



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**Assumption:**  $\ell$  is Lipschitz and  $H \in C^3(\mathbb{R})$  is uniform convex at  $\infty$ , that is

$$\lim_{x \rightarrow \pm\infty} H''(x) = c_{H,\pm} > 0.$$

Initial data  $\rho^0$  has finite second moment  $M_2(\rho^0) < \infty$   
and bounded free energy  $\mathcal{F}(\rho^0) < \infty$ .

**Well-posedness:** via direct method and convexity of  $\mathcal{F}$  and  $\mathcal{M}^{\ell}$

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### Passage to the limit $h \rightarrow 0$

Let  $\rho_h$  be a piecewise constant interpolation of  $\{\rho^k\}$  solution of the scheme, then  $\rho_h$  converges weakly to a weak solution of the constrained Fokker-Planck equation.

- Time-discrete approximation of the weak formulation: For all  $\zeta \in C_c^\infty(\mathbb{R})$

$$\left| \int_{\mathbb{R}} \frac{\rho_h^k - \rho_h^{k-1}}{h} \zeta + \int \left( (H'(x) - \sigma_h^k) \partial_x \zeta - \partial_{xx} \zeta \right) \rho_h^k \right| \leq \sup_{\mathbb{R}} \frac{|\partial_{xx} \zeta|}{2} \frac{1}{h} W_2^2(\rho_h^{k-1}, \rho_h^k),$$

with the discrete Lagrange multiplier  $\sigma_h^k$  given by

$$\sigma_h^k = \int_{\mathbb{R}} H'(x) \rho_h^k(x) dx + \frac{\ell(kh) - \ell((k-1)h)}{h}.$$

- A priori estimates on moments, entropy, energy and Wasserstein distance.
- Uniform convergence of the Lagrange multiplier via an Arzela-Ascoli argument

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# Convergence to equilibrium

### Naive application of entropy method fails:

- System is **not thermodynamically closed** for  $\dot{\ell} \neq 0$

$$\frac{d}{dt} \mathcal{F}(\rho(t)) = -\mathcal{D}(\rho(t), \sigma(t)) + \sigma(t) \dot{\ell}(t)$$

with

$$\mathcal{D}(\rho, \sigma) := \int |\nu^2 \partial_x \log \rho + H' - \sigma|^2 dx = \int \left| \nu^2 \partial_x \log \frac{\rho}{\gamma_\sigma} \right|^2 dx$$

and  $\gamma_\sigma(x) = \exp\left(-\frac{H(x) - \sigma x}{\nu^2}\right) / Z_\sigma$ .

- Naive use of logarithmic Sobolev inequality does not close the entropy relation, even for  $\dot{\ell} \equiv 0$

$$\frac{d}{dt} \mathcal{F}(\rho(t)) = -\mathcal{D}(\rho(t), \sigma(t)) \leq -\frac{\nu^2}{C_{\text{LSI}}} \mathcal{H}(\rho(t) | \gamma_{\sigma(t)})$$

Mismatch of reference state!

$$\frac{d}{dt} \mathcal{H}(\rho(t) | \gamma_0) \leq -\frac{1}{C_{\text{LSI}}} \mathcal{H}(\rho(t) | \gamma_{\sigma(t)})$$

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**Idea:** Constrained minimization of free energy

### Characterization of minimizer

There exists a unique minimizer of the constrained minimization problem

$$\arg \min_{\rho \in \mathcal{M}^\ell} \mathcal{F}(\rho) = \gamma_{\lambda(\ell)} \quad \text{with} \quad \gamma_{\lambda(\ell)}(x) = \frac{1}{Z_{\lambda(\ell)}} \exp\left(-\frac{H(x) - \lambda(\ell)x}{\nu^2}\right)$$

The function  $\lambda : \mathbb{R} \rightarrow \mathbb{R}$  implicitly defined by  $M_1(\gamma_{\lambda(\ell)}) = \ell$  satisfies the **Bi-Lipschitz estimate**

$$0 < c_{\text{var}} \leq \frac{dM_1(\gamma_\lambda)}{d\lambda} \leq C_{\text{var}} < \infty.$$

**Proof:**

- Convexity of  $\mathcal{F}$  and  $\mathcal{M}^\ell$  yields uniqueness
- Explicit construction for existence
- Bi-Lipschitz estimate:

$$\frac{dM_1(\gamma_\lambda)}{d\lambda} = \int x^2 \gamma_\lambda(x) dx - \left( \int x \gamma_\lambda(x) dx \right)^2 =: \text{var}(\gamma_\lambda).$$



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### Relative entropy magic

For all  $\eta \in \mathbb{R}$ , all  $\ell \in \mathbb{R}$  and all  $\rho \in \mathcal{M}^\ell$  it holds

$$\frac{C_{\text{var}}}{2}(\eta - \lambda(\ell))^2 \leq \mathcal{H}(\rho|\gamma_\eta) - \mathcal{H}(\rho|\gamma_{\lambda(\ell)}) \leq \frac{C_{\text{var}}}{2}(\eta - \lambda(\ell))^2$$

- Uniform bounds on moment, Lagrange multiplier and free energy

$$\sup_{t \geq 0} \max\{M_2(\rho(t)), \sigma(t), \mathcal{F}(\rho(t))\} \leq C.$$

- Uniform logarithmic Sobolev estimate

$$\forall |\sigma| \leq C : \quad \mathcal{H}(\rho|\gamma_\sigma) \leq C_{\text{LSI}} \mathcal{D}(\rho, \sigma).$$

The constant  $C_{\text{LSI}}^{-1}$  characterizes the rate of convergence to equilibrium

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### Convergence to equilibrium

Suppose the dynamical constraint becomes constants

$$\dot{\ell} \in L^1(\mathbb{R}_+) \quad \Rightarrow \quad \lim_{t \rightarrow \infty} \ell(t) = \ell^*.$$

Then

$$\mathcal{H}(\rho(t) | \gamma_{\lambda(\ell(t))}) \leq e^{-\frac{t}{C_{\text{LSI}}}} \mathcal{H}(\rho^0 | \gamma_{\lambda(\ell^0)}) + C_{\ell, \sigma} \int_0^t e^{-\frac{t-s}{C_{\text{LSI}}}} |\dot{\ell}(s)| ds.$$

with  $C_{\ell, \sigma} = \|\sigma\|_{\infty} + \|\lambda \circ \ell\|_{\infty}$ .

**Proof:**

$$\begin{aligned} \frac{d}{dt} \mathcal{H}(\rho(t) | \gamma_{\lambda(\ell(t))}) &= -\mathcal{D}(\rho(t), \sigma(t)) + \dot{\ell}(t)(\sigma(t) - \lambda(\ell(t))) \\ &\leq -\frac{1}{C_{\text{LSI}}} \mathcal{H}(\rho(t) | \gamma_{\sigma(t)}) + C_{\sigma, \ell} |\dot{\ell}(t)| \\ &\leq -\frac{1}{C_{\text{LSI}}} \mathcal{H}(\rho(t) | \gamma_{\lambda(\ell(t))}) + C_{\sigma, \ell} |\dot{\ell}(t)| \end{aligned}$$

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with  $C_{\ell, \sigma} = \|\sigma\|_{\infty} + \|\lambda \circ \ell\|_{\infty}$ .

**Proof:**

$$\begin{aligned} \frac{d}{dt} \mathcal{H}(\rho(t) | \gamma_{\lambda(\ell(t))}) &= -\mathcal{D}(\rho(t), \sigma(t)) + \dot{\ell}(t)(\sigma(t) - \lambda(\ell(t))) \\ &\leq -\frac{1}{C_{\text{LSI}}} \mathcal{H}(\rho(t) | \gamma_{\sigma(t)}) + C_{\sigma, \ell} |\dot{\ell}(t)| \\ &\leq -\frac{1}{C_{\text{LSI}}} \mathcal{H}(\rho(t) | \gamma_{\lambda(\ell(t))}) + C_{\sigma, \ell} |\dot{\ell}(t)| \end{aligned}$$

- Fix reference state: For any  $\ell^* \in \mathbb{R}$ , any  $\ell \in \mathbb{R}$  and all  $\rho \in \mathcal{M}^\ell$

$$\mathcal{H}(\rho|\gamma_{\lambda(\ell^*)}) \leq \mathcal{H}(\rho|\gamma_{\lambda(\ell)}) + \frac{C_{\text{var}}}{2c_{\text{var}}^2} |\ell^* - \ell|^2.$$

- Let  $\ell$  converge exponentially:  $|\dot{\ell}(t)| \leq L_0 e^{-\kappa t}$ , then

$$\mathcal{H}(\rho(t)|\gamma_{\lambda(\ell^*)}) \leq C_0 e^{-\tau t} \quad \text{with} \quad \tau = \min\{C_{\text{LSI}}^{-1}, \kappa\}.$$

- Classical Csiszár-Pinsker inequality implies  $L^1$ -convergence

$$\int |\rho(t, x) - \gamma_{\lambda(\ell^*)}(x)| dx \leq \tilde{C}_0 e^{-\frac{\tau}{2} t}.$$

- A weighted Csiszár-Pinsker inequality due to [Bolley, Villani 2005] yields

$$(\sigma(t) - \lambda(\ell^*))^2 \leq \tilde{C} e^{-\tau t}$$

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**Thank you for your attention!**